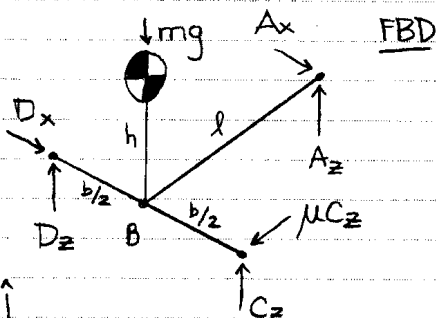


6.74 A branch gets caught in the right rear wheel of a tricycle causing the wheel to skid in the  $\hat{j}$ -direction. The wheels at A & D roll without slip and are massless. What's the acceleration of the tricycle?

- A & D roll w/o slip in  $\hat{j}$ -direction
- wheel at C skids in  $\hat{j}$ -direction



LMB:  $\Sigma \underline{F} = m\underline{a} = m a \hat{j}$

LMB ·  $\hat{j} \Rightarrow -\mu C_z = m a$  (1)

AMB<sub>B</sub> ·  $\hat{j} : \{ \Sigma \underline{M}_B = \underline{H}_B \} \cdot \hat{j}$  } moments about axis through B  
 $\Rightarrow -C_z(\frac{b}{2}) + D_z(\frac{b}{2}) = 0$  } // to  $\hat{j}$ -direction  
 $\therefore C_z = D_z$  (2)

AMB<sub>A</sub> ·  $\hat{i} : \{ \Sigma \underline{M}_A = \underline{H}_A \} \cdot \hat{i}$  } moments about axis through A  
 $\Rightarrow -C_z(l) - D_z(l) + mg(l) = -mah$  } // to  $\hat{i}$ -direction  
 $mg - (C_z + D_z) = -ma \frac{h}{l}$  (3)

(2) into (3):  $mg - 2C_z = -ma \frac{h}{l}$  (4)

(1) into (4):  $mg + \frac{2ma}{\mu} = -ma \frac{h}{l}$

$a(\frac{2}{\mu} + \frac{h}{l}) = -g$   
 $\Rightarrow a = \frac{-g}{\frac{2}{\mu} + \frac{h}{l}} = \frac{-\frac{\mu g}{2}}{1 + \frac{\mu h}{2l}}$

$\therefore \underline{a} = -\frac{\mu g}{2} \left( \frac{1}{1 + \frac{\mu h}{2l}} \right) \hat{j}$

Checks: ① When  $\mu = 0$ ,  $\underline{a} = 0$  (no forces in  $\hat{j}$ -direction  $\Rightarrow$  constant  $\underline{x}$ )

② When  $h = 0$ ,  $\underline{a} = -\frac{\mu g}{2} \hat{j}$   
 (Half of the rider's weight is supported at wheel C.)

position:  $\underline{r}(t) = R \cos(\dot{\theta}t) \hat{i} + R \sin(\dot{\theta}t) \hat{j}$

where  $\dot{\theta} = \text{constant}$

a)  $\underline{v} = \dot{\underline{r}} = R\dot{\theta}(-\sin(\dot{\theta}t)\hat{i} + \cos(\dot{\theta}t)\hat{j})$

$\therefore \underline{v} \neq \underline{0}$  since the velocity depends on time

b)  $\underline{v} \neq \text{constant}$  (see a)

c)  $|\underline{v}| = R\dot{\theta} = \text{constant}$  since  $\dot{\theta} = \text{const.}$

d)  $\underline{a} = \dot{\underline{v}} = -R\dot{\theta}^2(\cos(\dot{\theta}t)\hat{i} + \sin(\dot{\theta}t)\hat{j})$

$\therefore \underline{a} \neq \underline{0}$  since the acceleration depends on time

e)  $\underline{a} \neq \text{constant}$  (see d)

f)  $|\underline{a}| = R\dot{\theta}^2 = \text{constant}$  since  $\dot{\theta} = \text{const.}$

g)  $\underline{v} \cdot \underline{a} = -R^2\dot{\theta}^3(-\sin(\dot{\theta}t)\cos(\dot{\theta}t) + \cos(\dot{\theta}t)\sin(\dot{\theta}t)) = \underline{0}$

$\therefore \underline{v} \perp \underline{a}$

7.16

$\theta = \theta_0 \cos \lambda t$  ↙ given  
 $\Rightarrow \dot{\theta} = -\theta_0 \lambda \sin \lambda t$   
 $\ddot{\theta} = -\theta_0 \lambda^2 \cos \lambda t$

$\underline{a} = -R\dot{\theta}^2 \hat{e}_R + R\ddot{\theta} \hat{e}_\theta$

where  $\begin{aligned} \dot{\theta}^2 &= \theta_0^2 \lambda^2 \sin^2 \lambda t \\ &= \lambda^2 (\theta_0^2 - \theta_0^2 \cos^2 \lambda t) \\ &= \lambda^2 (\theta_0^2 - \theta^2) \end{aligned}$

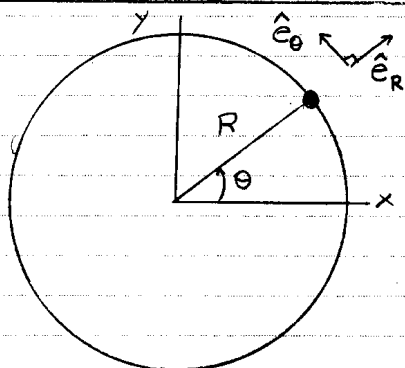
$\ddot{\theta} = -\lambda^2 \theta$

$\Rightarrow |\underline{a}|^2 = R^2 \dot{\theta}^4 + R^2 \ddot{\theta}^2$

$|\underline{a}|^2 = R^2 [\lambda^4 (\theta_0^2 - \theta^2)^2 + \lambda^4 \theta^2]$

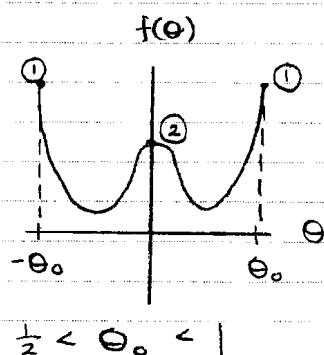
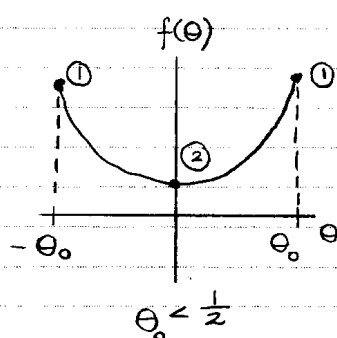
$|\underline{a}|^2(\theta) = R^2 \lambda^4 \underbrace{[(\theta_0^2 - \theta^2)^2 + \theta^2]}_{f(\theta)}$

(continued)



Maximum acceleration occurs at the maximum value of

$$f(\theta) = (\theta_0^2 - \theta^2)^2 + \theta^2 = \theta^4 + \theta^2(1 - 2\theta_0^2) + \theta_0^4$$

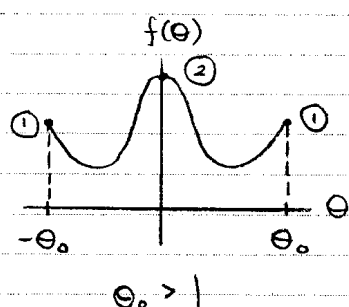


$$① f(\theta = \theta_0) = \theta_0^2$$

$$② f(\theta = 0) = \theta_0^4$$

∴ Max value of  $|a|$  will be at  $\theta = 0$  only when  $\theta_0 \geq 1$ .

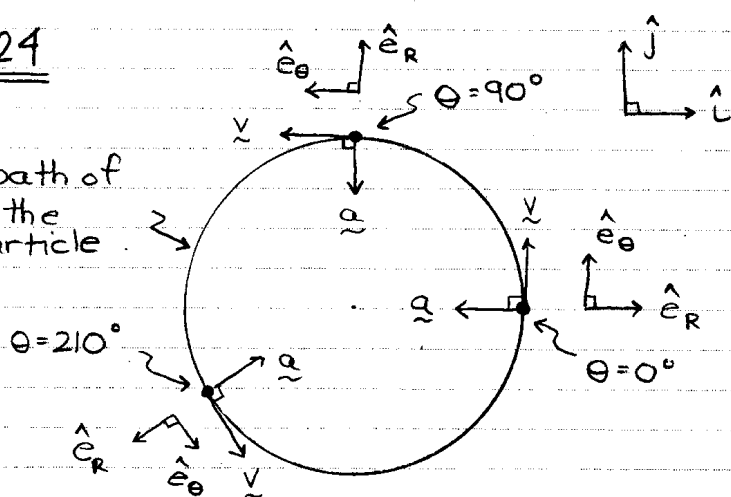
$$\Rightarrow |a| = R\dot{\theta}^2 \theta_0^2$$



## 7.24

a)

path of the particle



$$b) \dot{\theta} = \frac{2\pi}{T} = \text{constant}$$

$$\text{@ } \theta = 0^\circ : \hat{e}_R = \hat{i}, \hat{e}_\theta = \hat{j}$$

$$\underline{v} = R\dot{\theta}\hat{j}, \underline{a} = -R\dot{\theta}^2\hat{i}$$

$$\text{@ } \theta = 90^\circ : \hat{e}_R = \hat{j}, \hat{e}_\theta = -\hat{i}$$

$$\underline{v} = -R\dot{\theta}\hat{i}, \underline{a} = -R\dot{\theta}^2\hat{j}$$

(continued)

$$\textcircled{\theta=210^\circ}: \hat{e}_R = -\frac{\sqrt{3}}{2}\hat{i} - \frac{1}{2}\hat{j}$$

$$\hat{e}_\theta = \frac{1}{2}\hat{i} - \frac{\sqrt{3}}{2}\hat{j}$$

$$\underline{v} = R\dot{\theta} \left( \frac{1}{2}\hat{i} - \frac{\sqrt{3}}{2}\hat{j} \right)$$

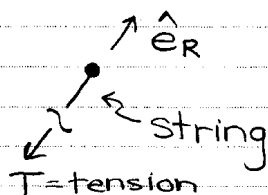
$$\underline{a} = R\dot{\theta}^2 \left( \frac{\sqrt{3}}{2}\hat{i} + \frac{1}{2}\hat{j} \right)$$

c) LMB:  $\Sigma \underline{F} = m\underline{a}$

FBD

$$\{-T\hat{e}_R = m(-R\dot{\theta}^2\hat{e}_R)\}$$

$$\{\} \cdot \hat{e}_R \Rightarrow \boxed{T = mR\dot{\theta}^2}$$



d) Since  $\dot{\theta} = \text{constant}$ , there is no tangential acceleration. Therefore, the radial tension is all that's needed to keep the motion going.

e)  $\underline{T} = -T\hat{e}_R = T\left(\frac{\sqrt{3}}{2}\hat{i} + \frac{1}{2}\hat{j}\right) @ \theta = 210^\circ$

$$\therefore \boxed{F_x = \frac{\sqrt{3}}{2}T = \frac{\sqrt{3}}{2}mR\dot{\theta}^2}$$

7.29

LMB:  $\Sigma \underline{F} = m\underline{a}$

$$\{-k(r-r_0)\hat{e}_R = m(-r\omega_0^2\hat{e}_R)\}$$

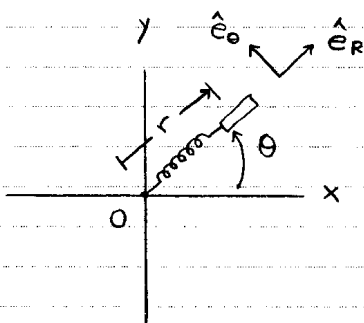
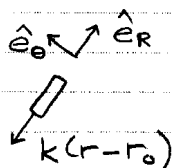
$$\{\} \cdot \hat{e}_R \Rightarrow -k(r-r_0) = -mr\omega_0^2$$

$$-kr + kr_0 = -mr\omega_0^2$$

$$r(k - m\omega_0^2) = kr_0$$

$$\therefore \boxed{r = \frac{kr_0}{k - m\omega_0^2}}$$

FBD



a) LMB:  $\Sigma \underline{F} = m \underline{a}$

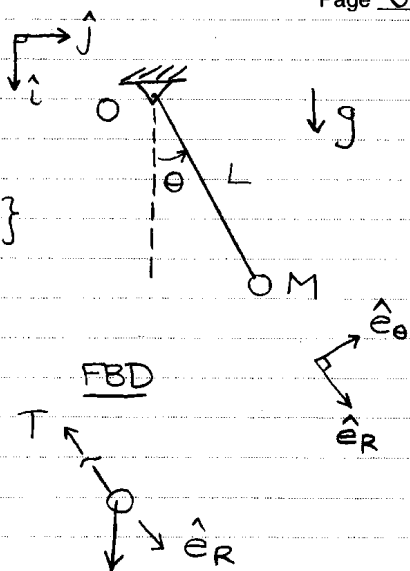
$$\{-T \hat{e}_R + mg \hat{i} = m(-L\ddot{\theta} \hat{e}_R + L\ddot{\theta} \hat{e}_\theta)\}$$

$$\{\} \cdot \hat{e}_\theta$$

$$\Rightarrow mg(\hat{i} \cdot \hat{e}_\theta) = mL\ddot{\theta}$$

where

$$\begin{aligned} \hat{i} \cdot \hat{e}_\theta &= \hat{i} \cdot (-\sin\theta \hat{i} + \cos\theta \hat{j}) \\ &= -\sin\theta \end{aligned}$$



$$\therefore -mg \sin\theta = mL\ddot{\theta}$$

$$\Rightarrow \ddot{\theta} = -\frac{g}{L} \sin\theta$$

b) LMB  $\cdot \hat{e}_R \Rightarrow -T + mg(\hat{i} \cdot \hat{e}_R) = -mL\dot{\theta}^2$

where  $\hat{i} \cdot \hat{e}_R = \hat{i} \cdot (\cos\theta \hat{i} + \sin\theta \hat{j}) = \cos\theta$

$$\Rightarrow T = mg \cos\theta + mL\dot{\theta}^2$$

c) The force at the hinge support is simply  $\underline{F}_0 = -T \hat{e}_R$  because the string is massless.

$$\underline{F}_0 = -T \hat{e}_R = -T(\cos\theta \hat{i} + \sin\theta \hat{j})$$

$$\Rightarrow \begin{cases} \underline{F}_0 \cdot \hat{i} = -T \cos\theta \\ \underline{F}_0 \cdot \hat{j} = -T \sin\theta \end{cases}$$

where T is given above

d) Let  $\alpha = \dot{\theta} \Rightarrow \dot{\alpha} = \ddot{\theta} = -\frac{g}{L} \sin\theta$

$$\Rightarrow \begin{cases} \dot{\theta} = \alpha \\ \dot{\alpha} = -\frac{g}{L} \sin\theta \end{cases}$$

e) see next page

(continued)

f) Max. tension = 29.9999... N

g)  $T = 2.3452$  s  $\leftarrow$  get by fine tuning tspan, interpolation, or event detection

global g L M

g=10; L=1; M=1;

tspan=[0 5];

theta0=pi/2; thetadot0=0;

z0=[theta0 thetadot0];

options=odeset('AbsTol',1e-12,'RelTol',1e-12);

[t,z]=ode45('deriv43',tspan,z0,options);

theta=z(:,1); thetadot=z(:,2);

tension=M\*g\*cos(theta)+M\*L\*thetadot.^2;

max(tension)

plot(t,theta)

hold on

plot(t,thetadot,'--')

plot(t,tension,':')

function zdot=deriv43(t,z)

global g L M

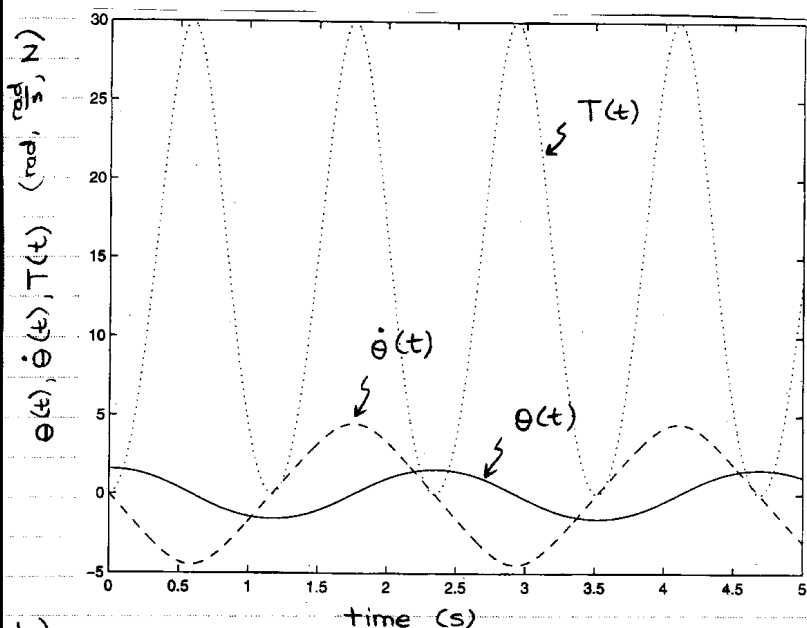
theta=z(1); alpha=z(2);

thetadot=alpha;

alphadot=(-g/L)\*sin(theta);

zdot=[thetadot alphadot]';

Plots for 7.43



h)

Note: ① Maximum tension is independent of string length, ② The max. tension can be made closer to 30 N by decreasing the integration tolerance, ③ Using  $\sin \theta \approx \theta$  get  $T = 1.9869$  s (cf.  $T = 2.3452$  s above).

7.44

$$T = \frac{2\pi}{\lambda} = 2\pi\sqrt{\frac{L}{g}}$$

LMB:  $\Sigma \underline{F} = m\underline{a}$ 

$$\{-N\hat{e}_R - \mu N\hat{e}_\theta = m(-\frac{v^2}{R}\hat{e}_R + \dot{v}\hat{e}_\theta)\}$$

$$\{\} \cdot \hat{e}_R \Rightarrow N = m\frac{v^2}{R} \quad (1)$$

$$\{\} \cdot \hat{e}_\theta \Rightarrow -\mu N = m\dot{v} \quad (2)$$

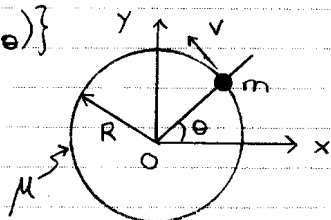
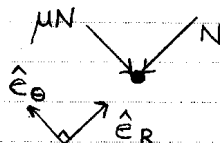
$$(1) \text{ into } (2): \quad \dot{v} = -\frac{\mu v^2}{R}$$

Since  $v = R\dot{\theta}$ ,

$$\frac{dv}{dt} = -\mu v \frac{d\theta}{dt}$$

$$\Rightarrow \frac{dv}{d\theta} = -\mu v$$

$$\therefore v(\theta) = v_0 e^{-\mu\theta}$$

FBDused  $v(0) = v_0$