

A turntable oscillates with displacement $\vec{x}(t) = A\sin(\omega t)\hat{i}$. A bug walks along DE with v_0 relative to the turntable.

At the moment, $x=A$, DE is parallel to z and B . $|CB|=a$. v_0 is given. find \vec{a}_b .

We consider a frame B attached to the turntable with its origin at C .

$$\vec{a}_b = \vec{a}_{b'} + \vec{a}_{\text{rel}} + \vec{a}_{\text{rel}}$$

where $\vec{a}_{b'} =$ acceleration of a point B' that was fixed in E and at the moment coincides with B .

$$\vec{a}_{\text{rel}} = \text{Coriolis acceleration}$$

$\vec{a}_{\text{rel}} =$ Acceleration of B relative to frame B

$$\vec{a}_b = \vec{a}_{b'} + \vec{a}_B \times \vec{r}_{B/C} + \vec{\omega}_B \times (\vec{a}_B \times \vec{r}_{B/C})$$

$$= -A\omega^2 \sin(\omega t)\hat{i} + \omega_B \hat{j} \times \hat{a}_B \hat{i} + \omega_B \hat{j} \times (\omega_B \hat{j} \times \hat{a}_B \hat{i})$$

$$= -(A\omega^2 + \omega_B^2)\hat{i} - a_{B,k}\hat{k} \quad (\text{since at this instant } A = A\sin(\omega t), \therefore \sin(\omega t) = 1)$$

$$\vec{a}_{\text{rel}} = 2\vec{\omega}_B \times \vec{a}_{\text{rel}}$$

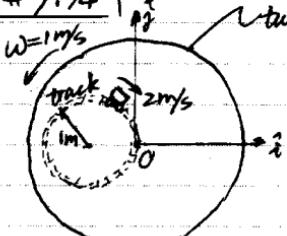
$$= 2\omega_B \hat{j} \times (-v_0) \hat{k}$$

$$= -2\omega_B v_0 \hat{i}$$

$$\vec{a}_{\text{rel}} = 0$$

$$\therefore \vec{a}_b = -(A\omega^2 + \omega_B^2 + 2\omega_B v_0)\hat{i} - a_{B,k}\hat{k} \text{ m/s}^2$$

9.74



0.1 kg toy train is going clockwise at a constant rate.

turntable is rotating counterclockwise.

At the instant, train is pointing due south ($-\hat{j}$) and is at the center O of the turntable.

Consider the frame B attached to the turntable, origin at (x, y)

a) velocity of the train relative to the turntable?

$$\vec{v}_{B/B} = -2\hat{j} \text{ m/s}$$

#

#9.74 Cont'd

b) What's the absolute velocity?

$$\vec{v} = \vec{v}_{B'} + \vec{v}_{rel}$$

B' is the point coinciding with the train but fixed on the rotating frame \mathcal{D} . at the instant.

$$\vec{v}_{B'} = \vec{v}_0 + \vec{\omega}_0 \times \vec{r}_{B'/0} = 0 \quad \vec{v}_{rel} = -2\hat{j} \text{ m/s}$$

$$\therefore \vec{v} = -2\hat{j} \text{ m/s}$$

c) Acceleration relative to the turntable?

$$\vec{a}_{B'} = -\frac{1}{r}\hat{i} = -4\hat{i} \text{ m/s}^2$$

d) Absolute acceleration?

$$\vec{a} = \vec{a}_{B'} + \vec{a}_{rot} + \vec{a}_{rel}$$

B' is the point coinciding with the train but fixed on the rotating frame \mathcal{D} at the instant

$$\text{where } \vec{a}_{B'} = \vec{v}_0 + \vec{\omega}_0 \times \vec{r}_{B'/0} + \vec{\omega}_0 \times (\vec{\omega}_0 \times \vec{r}_{B'/0}) = 0$$

$$\vec{a}_{rot} = 2\vec{\omega}_0 \times \vec{v}_{rel} = 2(1\hat{k}) \times (-2\hat{j}) = 4\hat{i} \text{ m/s}^2$$

$$\vec{a}_{rel} = -4\hat{i} \text{ m/s}^2$$

$$\therefore \vec{a} = 0 \text{ m/s}^2$$

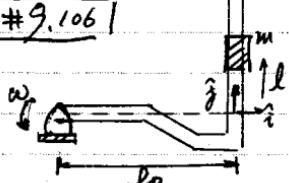
e) Total force acting on train?

$$2\vec{F} = m\vec{a} = 0 \quad \therefore \text{The total force is 0.}$$

f) Sketch.

Suppose $t=0$, the train is at the center.

goes back & forth along the vertical line.

#9.106

No friction, bead starts at $l=0$, $\dot{l}=v_0$, $\omega=w_0$ constant.
Find $l(t)$

Set up the frame $\mathcal{D}(x-y)$ attached the instantaneously vertical body rotating with

FBD:



$$\vec{a} = \vec{a}_g + \vec{a}_{rot} + \vec{a}_{rel}$$

$$\vec{a}_g = \vec{a}_0 + \vec{\omega}_0 \times \vec{r}_{m/g} + \vec{\omega}_0 \times (\vec{\omega}_0 \times \vec{r}_{m/g})$$

$$= -l_0 w_0^2 \hat{i} - l w_0^2 \hat{j}$$

$$\vec{a}_{rot} = 2\vec{\omega}_0 \times \vec{v}_{rel} = -2l w_0 \hat{i}$$

$$\vec{a}_{rel} = \dot{l} \hat{j}$$

$$\therefore \vec{a} = (-l\omega_0^2 - 2i\omega_0)\hat{i} + (-l\omega_0^2 + l)\hat{j}$$

$$\sum \vec{F} = m\vec{a}$$

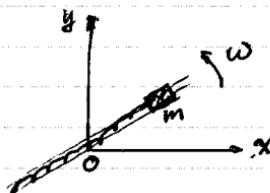
$$\{ \} \cdot \hat{j}: l - l\omega_0^2 = 0$$

$$\therefore l = A_1 e^{i\omega_0 t} + A_2 e^{-i\omega_0 t}$$

$$t=0, l=0, i=\omega_0 v_0 \quad \therefore \begin{cases} A_1 + A_2 = 0 \\ A_1 \omega_0 - A_2 \omega_0 = v_0 \end{cases} \quad \therefore \begin{cases} A_1 = \frac{v_0}{2\omega_0} \\ A_2 = -\frac{v_0}{2\omega_0} \end{cases}$$

$$\therefore l(t) = \frac{v_0}{2\omega_0} (e^{i\omega_0 t} - e^{-i\omega_0 t})$$

9.113

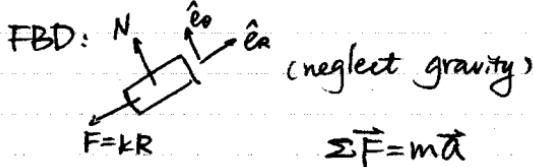


No friction. ω constant.
Spring constant k .

Spring is relaxed when bead is at the center position.

Assume $t=0$, $R=d$, $\dot{R}=0$. If needed, may assume $k > m\omega^2$

a) Derive eqn for the position of the bead $R(t)$.



$$\therefore -kR\hat{e}_R + N\hat{e}_\theta = m[(\ddot{R} - R\dot{\theta}^2)\hat{e}_R + (2R\dot{\theta} + R\ddot{\theta})\hat{e}_\theta]$$

$$\{ \} \cdot \hat{e}_R: -kR = m\ddot{R} - mR\dot{\theta}^2 \quad (*)$$

$$\{ \} \cdot \hat{e}_\theta: N = 2mR\dot{\theta}$$

$$\text{From } (*): \ddot{R} + \frac{k-m\omega^2}{m}R = 0 \quad (\text{assume } k > m\omega^2)$$

$$\therefore R = A_1 \sin \sqrt{\frac{k-m\omega^2}{m}} t + A_2 \cos \sqrt{\frac{k-m\omega^2}{m}} t$$

$$\text{small } t=0, R=d, \dot{R}=0 \quad \therefore A_1=0, A_2=d.$$

$$\therefore R = d \cos \sqrt{\frac{k-m\omega^2}{m}} t$$

b) How is the motion affected by large k vs small k ?

from $R = d \cos \sqrt{\frac{k-m\omega^2}{m}} t$, we can see that when k is very large, $\sqrt{\frac{k-m\omega^2}{m}}$ is large. (when $k \gg m\omega^2$, $\sqrt{\frac{k-m\omega^2}{m}} \approx \sqrt{k}$ const) the frequency of the motion is also large.

When k is small, frequency decreases.

c) When bead passes through the center position, $R=0$

$$\therefore \cos\sqrt{\frac{k-m\omega^2}{m}}t = 0 \quad \sin\sqrt{\frac{k-m\omega^2}{m}}t = 1$$

$$\therefore N = 2mR\ddot{\theta} = 2md\sqrt{\frac{k-m\omega^2}{m}}(-\sin\sqrt{\frac{k-m\omega^2}{m}}t)w$$

$$\therefore |N| = 2md\omega\sqrt{\frac{k-m\omega^2}{m}}$$

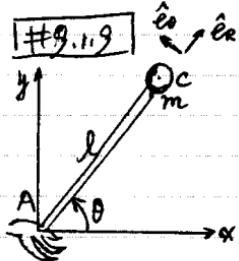
the magnitude of the force acting on the bead.

d) Coriolis acceleration:

$$2R\dot{\theta}\hat{e}_0 = -2md\omega\sqrt{\frac{k-m\omega^2}{m}}\sin\sqrt{\frac{k-m\omega^2}{m}}t\hat{e}_0$$

Example: you throw a particle and find it falls to ground a distance east of the point supposed to be.

$\begin{matrix} v \\ \downarrow \\ \omega \end{matrix} \Rightarrow \begin{matrix} \hat{e}_0 \\ \downarrow \\ \omega \end{matrix}$ Due to the rotation of the earth.
 $\vec{a} = 2\vec{\omega} \times \vec{v}_{\text{rel.}}$ Coriolis acceleration, eastward.

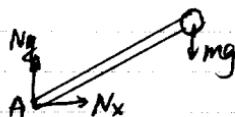


rod on palm A, length l, mass $m\omega l^2$

Assume $\dot{\theta}$ & θ are known at the instant of interest. hand is accelerating

$$\vec{a}_{\text{hand}} = a_{hx}\hat{i} + a_{hy}\hat{j}$$

a) Draw FBD



b) Assume hand is stationary. Solve for $\ddot{\theta}$

From the acceleration in polar coordinates:

$$(R-R\dot{\theta}^2)\hat{e}_R + (2R\dot{\theta}+R\ddot{\theta})\hat{e}_{\theta}$$

$$R=l \quad \therefore R=\ddot{R}=0 \quad \therefore \vec{a}_c = -R\dot{\theta}^2\hat{e}_R + R\ddot{\theta}\hat{e}_{\theta}$$

$$\vec{e}_R = \cos\theta\hat{i} + \sin\theta\hat{j} \quad \vec{e}_{\theta} = -\sin\theta\hat{i} + \cos\theta\hat{j}$$

$$\therefore \vec{F}_A = \vec{r}_A \times m\vec{a}_c$$

$$= l(\cos\theta\hat{i} + \sin\theta\hat{j}) \times m[-l\dot{\theta}^2(\cos\theta\hat{i} + \sin\theta\hat{j}) + l\ddot{\theta}(-\sin\theta\hat{i} + \cos\theta\hat{j})] \quad \text{in cartesian}$$

$$= ml^2\ddot{\theta}\hat{k} \quad \text{in polar coordinates}$$

$$\sum \vec{M}_A = \vec{r}_A \times (-mg\hat{j}) = -mgl\cos\theta\hat{k}$$

$$\therefore \int \vec{b} \cdot \hat{k}: ml^2\ddot{\theta} = -mgl\cos\theta$$

$$\therefore \ddot{\theta} = -\frac{g\cos\theta}{l}$$

c) If hand is not stationary, but $\dot{\theta}$ is given.

$$\sum \vec{F} = m\vec{a} \quad \text{where } \vec{a} = \vec{a}_{\text{hand}} + \vec{a}_{\text{chand}}$$

$$\{ \cdot \hat{j}: \quad N_y - mg = ma_{cy} + m\ddot{a}_{hy}$$

$$\text{And } a_{cy} = -l\ddot{\theta}\sin\theta + l\ddot{\theta}\cos\theta.$$

$$\therefore N_y = mg + m\ddot{a}_{hy} + m(l\ddot{\theta}\cos\theta - l\ddot{\theta}\sin\theta)$$