

A turntable oscillates with displacement $\tilde{x}(t) = A \sin(\pi t)$. A bug walks along DE with v_b relative to the turntable.

At the moment, $x=A$, DE is parallel to $\hat{y}$ axis. $\omega = \omega_0$ given. find $\tilde{a}_B$ to $\hat{z}$ axis.

We consider a frame B attached to the turntable with its origin at C.

$$\tilde{a}_B = \tilde{a}_C + \tilde{a}_{rel} + \tilde{a}_{Cor}$$

where $\tilde{a}_B$ = acceleration of a point B' that is fixed in B and at the moment coincides with B.

$\tilde{a}_{rel}$ = Coriolis acceleration

$\tilde{a}_{rel} = \text{Acceleration of B relative to frame B}$

$$\tilde{a}_B = \tilde{a}_C + \omega_B \times \tilde{r}_{BC} + \tilde{a}_{rel} + \omega_B \times (\omega_B \times \tilde{r}_{BC})$$

$$= -A\pi^2 \sin(\pi t) \hat{i} + \omega_B \hat{z} \times A\pi^2 \hat{i} + \omega_B \hat{z} \times (\omega_B \hat{z} \times A\pi^2 \hat{i})$$

$$= -(A\pi^2 + \omega_B^2 A) \hat{i} - \omega_B \hat{k} \quad (\text{since at this instant } A = A \sin(\pi t) \therefore \sin(\pi t) = 1)$$

$$\tilde{a}_{rel} = 2\omega_B \times \tilde{v}_{rel}$$

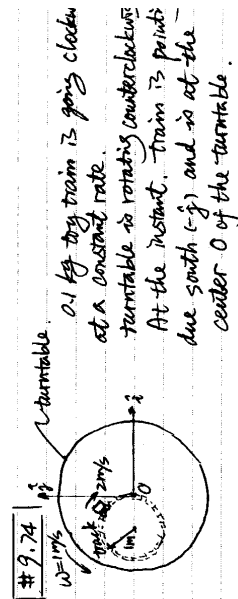
$$= 2\omega_B \hat{z} \times (-v_b) \hat{i}$$

$$= -2\omega_B v_b \hat{j}$$

$$\tilde{a}_{rel} = 0$$

$$\therefore \tilde{a}_B = -(A\pi^2 + \omega_B^2 A) \hat{i} - \omega_B \hat{k} \quad \text{m/s}^2$$

#9.74



Consider the frame B attached to the turntable, origin at (x,y).

a) velocity of the train relative to the turntable?

$$\tilde{v}_B = -2\hat{j} \text{ m/s}$$

b) What's the absolute velocity?

$\tilde{v} = \tilde{v}_B + \tilde{v}_{rel}$ B' is the point coinciding with the train but fixed on the rotating frame B at the instant.

$$\tilde{v}_B' = \tilde{v}_B + \omega_B \times \tilde{r}_{BO} = 0 \quad \tilde{v}_{rel} = -2\hat{j} \text{ m/s}$$

$$\therefore \tilde{v} = -2\hat{j} \text{ m/s}$$

c) Acceleration relative to the turntable?

$$\tilde{a}_B = -\frac{v}{r} \hat{i} = -4\hat{i} \text{ m/s}^2$$

d) Absolute acceleration?

$\tilde{a} = \tilde{a}_B + \tilde{a}_{rel} + \tilde{a}_{Cor}$ B' is the point coinciding with the train but fixed on the rotating frame B at the instant.

$$\text{where } \tilde{a}_B = \tilde{a}_B + \omega_B \times \tilde{r}_{BO} + \omega_B \times (\omega_B \times \tilde{r}_{BO}) = 0$$

$$\tilde{a}_{rel} = 2\omega_B \times \tilde{v}_{rel} = 2(\hat{k}) \times (-2\hat{j}) = 4\hat{i} \text{ m/s}^2$$

$$\tilde{a}_{Cor} = -4\hat{i} \text{ m/s}^2$$

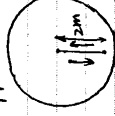
$$\therefore \tilde{a} = 0 \text{ m/s}^2$$

e) Total force acting on train?

$$\Sigma \tilde{F} = m\tilde{a} = 0 \quad \therefore \text{The total force is 0.}$$

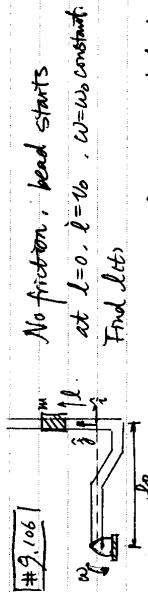
f) Sketch.

suppose $t=0$, the train is at the center.



goes back & forth along the vertical line.

#9.106



Set up the frame B(x,y) attached to the instantaneously vertical body rotating with

FBD:

$$\tilde{a} = \tilde{a}_B + \tilde{a}_{rel} + \tilde{a}_{Cor}$$

$$\tilde{a}_B = \tilde{a}_B + \omega_B \times \tilde{r}_{BO} + \omega_B \times (\omega_B \times \tilde{r}_{BO})$$

$$= -\omega_B^2 \tilde{r}_{BO} = -\omega_B^2 l \hat{i}$$

$$\tilde{a}_{rel} = 2\omega_B \times \tilde{v}_{rel} = -2\omega_B \hat{i} \dot{l}$$

$$\tilde{a}_{Cor} = \dot{l} \hat{j}$$

$$\therefore \tilde{a} = (-\omega_B^2 l - 2\omega_B \dot{l}) \hat{i} + (-\omega_B^2 \dot{l}) \hat{j}$$

$$\Sigma \tilde{F} = m\tilde{a}$$

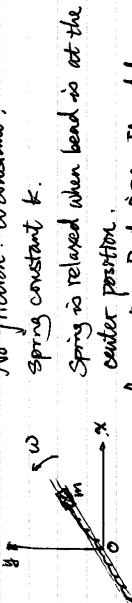
$$\left\{ \begin{array}{l} \hat{i} \cdot \hat{j} : \ddot{l} - \omega_B^2 l = 0 \\ \hat{j} \cdot \hat{j} : \ddot{l} - \omega_B^2 \dot{l} = 0 \end{array} \right.$$

$$\therefore \ddot{l} = A_1 e^{\omega_B t} + A_2 e^{-\omega_B t}$$

$$t=0, l=0, \dot{l} = \omega_B l_0 \quad \therefore \left\{ \begin{array}{l} A_1 + A_2 = 0 \\ A_1 \omega_B - A_2 \omega_B = \omega_B \end{array} \right. \quad \therefore \left\{ \begin{array}{l} A_1 = \frac{l_0}{2} \\ A_2 = -\frac{l_0}{2} \end{array} \right.$$

$$\therefore l(t) = \frac{l_0}{2\omega_B} (e^{\omega_B t} - e^{-\omega_B t})$$

#9.113



No friction. ω constant.

Spring constant k.

Spring is relaxed when bead is at the center position.

Assume $t=0, R=d, \dot{R}=0$. If needed, may assume $k > m\omega^2$.

a) Derive eqn for the position of the bead R(t).

FBD: $N = \hat{e}_\theta$ (neglect gravity)

$$\Sigma \tilde{F} = m\tilde{a}$$

$$\therefore -kR + N\hat{e}_\theta = m[(\ddot{R} - R\dot{\theta}^2)\hat{e}_r + (2\dot{R}\dot{\theta} + R\ddot{\theta})\hat{e}_\theta]$$

$$\left\{ \begin{array}{l} \hat{e}_r : -kR = m\ddot{R} - mR\dot{\theta}^2 \\ \hat{e}_\theta : N = 2mR\dot{\theta} \end{array} \right.$$

From (a): $\ddot{R} + \frac{k}{m}R = 0$ (assume $k > m\omega^2$)

$$\therefore R = A \sin \sqrt{\frac{k-m\omega^2}{m}} t + A_2 \cos \sqrt{\frac{k-m\omega^2}{m}} t$$

$$\text{since } t=0, R=d, \dot{R}=0 \quad \therefore A_1=0, A_2=d.$$

$$\therefore R = d \cos \sqrt{\frac{k-m\omega^2}{m}} t$$

b) How is the motion affected by large k vs small k?

from $R = d \cos \sqrt{\frac{k-m\omega^2}{m}} t$, we can see that when k is very large, $\sqrt{\frac{k-m\omega^2}{m}}$ is large. (when $k \gg m\omega^2$, $\sqrt{\frac{k-m\omega^2}{m}} \approx \sqrt{\frac{k}{m}}$)

the frequency of the motion is also large.

When k is small, frequency decreases.

c) When bead passes through the center position, $R=0$

$$\therefore \cos \sqrt{\frac{k-m\omega^2}{m}} t = 0 \quad \sin \sqrt{\frac{k-m\omega^2}{m}} t = 1$$

$$\therefore N = 2mR\ddot{\theta} = 2md \sqrt{\frac{k-m\omega^2}{m}} (-\sin \sqrt{\frac{k-m\omega^2}{m}} t) \omega$$

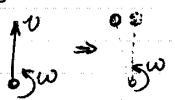
$$\therefore |N| = 2md\omega \sqrt{\frac{k-m\omega^2}{m}}$$

the magnitude of the force acting on the bead.

d) Coriolis acceleration:

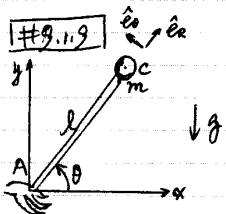
$$2\dot{R}\dot{\theta}\hat{e}_\theta = -2md\omega \sqrt{\frac{k-m\omega^2}{m}} \sin \sqrt{\frac{k-m\omega^2}{m}} t \hat{e}_\theta$$

Example: you throw a particle and find it falls to ground a distance east of the point supposed to be.



Due to the rotation of the earth.

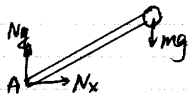
$\vec{a} = 2\vec{\omega} \times \vec{v}_{rel}$. Coriolis acceleration, eastward.



rod on palm A, length l , mass m . Assume $\dot{\theta}$ and $\ddot{\theta}$ are known at the instant of interest. hand is accelerating

$$\vec{a}_{hand} = a_{hx}\hat{i} + a_{hy}\hat{j}$$

a) Draw FBD



b) Assume hand is stationary. Solve for $\ddot{\theta}$

From the acceleration in polar coordinates:

$$(\ddot{R} - R\dot{\theta}^2)\hat{e}_r + (2\dot{R}\dot{\theta} + R\ddot{\theta})\hat{e}_\theta$$

$$R=l \therefore \ddot{R}=\dot{R}=0 \quad \therefore \vec{a}_c = -R\dot{\theta}^2\hat{e}_r + R\ddot{\theta}\hat{e}_\theta$$

$$\hat{e}_r = \cos\theta\hat{i} + \sin\theta\hat{j} \quad \hat{e}_\theta = -\sin\theta\hat{i} + \cos\theta\hat{j}$$

$$\therefore \vec{\tau}_A = \vec{r}_A \times m\vec{a}_c$$

$$= l(\cos\theta\hat{i} + \sin\theta\hat{j}) \times m[-l\dot{\theta}^2(\cos\theta\hat{i} + \sin\theta\hat{j}) + l\ddot{\theta}(-\sin\theta\hat{i} + \cos\theta\hat{j})] \quad \text{in cartesian}$$

$$= ml^2\ddot{\theta}\hat{k} \quad \text{in polar coordinates}$$

$$\Sigma \vec{\tau}_A = \vec{r}_A \times (-mg\hat{j}) = -mgl\cos\theta\hat{k}$$

$$\therefore \int \hat{k} : m\ddot{\theta} = -mgl\cos\theta$$

$$\therefore \ddot{\theta} = -\frac{g\cos\theta}{l}$$

c) If hand is not stationary, but $\dot{\theta}$ is given.

$$\Sigma \vec{F} = m\vec{a} \quad \text{where } \vec{a} = \vec{a}_{hand} + \vec{a}_{g/hand}$$

$$\hat{i} \cdot \hat{j} : N_y - mg = ma_{cy} + ma_{hy}$$

$$\text{And } a_{cy} = -l\dot{\theta}^2\sin\theta + l\ddot{\theta}\cos\theta$$

$$\therefore N_y = mg + ma_{hy} + m(l\ddot{\theta}\cos\theta - l\dot{\theta}^2\sin\theta)$$